

# The ACT Event Law: A Candidate Formulation

One Postulate, Four Theorems

A Working Note

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## Abstract

Every reviewer of ACT, internal and external, has converged on one missing object: an explicit event law—a hazard  $\lambda_k(t|\rho, E)$  for the occurrence of event  $k$  and a map  $\rho \rightarrow \mathcal{J}_k[\rho]$  stating what an event physically does. This note supplies a candidate formulation. ACT’s event dynamics is proposed as a piecewise-deterministic process (PDP): between events, the ACT-modified Schrödinger evolution of the Mathematical Supplement; events occur in the interaction-selected pointer basis at hazard  $\lambda_k(t) = \Lambda(t) p_k(t)$ , where  $p_k$  is the pointer-basis population and  $\Lambda(t) = d\Phi/dt$  is the rate of irreversible record formation already defined by the anchoring functional; the event map is the projective update  $\mathcal{J}_k[\rho] = P_k \rho P_k / p_k$ . The formulation’s value is what becomes *provable* once it is written down. Four theorems, each verified numerically to machine precision: (1) **Born uniqueness**—within this event class, the linear hazard  $\lambda_k \propto p_k$  is the *only* choice whose unconditioned dynamics is linear; any nonlinear hazard enables superluminal signalling through entangled pairs. The Born weights are therefore not assumed but forced by no-signalling. (2) **No-signalling**—local event maps leave spacelike-separated marginals exactly invariant. (3) **Order-invariance**—commuting spacelike event maps yield joint outcome statistics independent of event ordering, so observed probabilities are frame-independent even though the model’s event sequence is not. (4) **Ensemble consistency**—averaging the PDP reproduces the Joos–Zeh dephasing master equation exactly, so the law cannot conflict with any ensemble-level quantum prediction. Energy accounting is stated honestly: conserved in expectation up to a calculable per-event coherence-energy offset drawn against the bath. A structural bonus emerges from the law itself: since  $\Lambda = d\Phi/dt$ , the survival probability of the pre-event state is  $S(t) = e^{-\Phi(t)}$ —so  $\Phi$  is the cumulative event hazard,  $\Phi \sim 1$  is the e-folding scale rather than a boundary, and the long-criticized “arbitrary threshold” dissolves: no threshold is postulated at all. The note closes with the ledger: what this formulation postulates, what it derives, and what remains open—a fully covariant field-theoretic version chief among them.

## 1 What Was Missing

The [Mathematical Supplement](#) derives record formation: the anchoring functional  $\Phi_{\mathcal{O}}$  grows, off-diagonal coherences decay, and pointer components stabilize. It then attaches single-outcome realization to threshold crossing by an explicitly labeled bridge postulate. Five rounds of review identified the same gap: the postulate names *that* one outcome is realized but supplies no law for *which, when, with what probability, or with what physical consequence*. The required object is a stochastic event law:

$$\lambda_k(t | \rho, E) \quad \text{and} \quad \rho \longrightarrow \mathcal{J}_k[\rho]. \quad (1)$$

This note proposes one, in the most conservative mathematical class available—piecewise-deterministic processes, the same class that underlies quantum-jump trajectory theory (Breuer–Petruccione) and GRW hits—and then proves what the proposal buys.

## 2 The Event Law

**Postulate 1** (ACT Event Law). *Let  $\{P_k\}$  be the interaction-selected pointer projectors (einselection) for a system whose anchoring functional  $\Phi(t)$  is accumulating through coupling to environment  $E$ , and let  $p_k(t) = \text{Tr}[P_k \rho_S(t)]$ . The physical history of the system is a piecewise-deterministic process:*

1. **Flow:** *between events,  $\rho_S$  evolves by the ACT-modified dynamics of the Supplement (Schwinger–Keldysh reduced evolution; ordinary Schrödinger flow in the decoupled limit).*
2. **Hazard:** *event  $k$  occurs at instantaneous rate*

$$\lambda_k(t) = \Lambda(t) p_k(t), \quad \Lambda(t) \equiv \left. \frac{d\Phi}{dt} \right|_{irr}, \quad (2)$$

*where  $\Lambda$  is the rate of irreversible record formation: only environmental entanglement that has itself crossed the anchoring criterion contributes (reversible which-path correlations, as in quantum-eraser configurations, contribute zero hazard).*

3. **Event:** *when event  $k$  fires,*

$$\rho_S \longrightarrow \mathcal{J}_k[\rho_S] = \frac{P_k \rho_S P_k}{p_k}, \quad (3)$$

*applied to the system-plus-local-record degrees of freedom. The realized component enters the causal order (temporal existence); the unrealized components become causally inert: they enter no future hazard, source no future coupling, and constitute unrealized spectral possibility in the sense already used by the manuscript.*

Three remarks before the theorems. First, this is one postulate, not three: the PDP class, the pointer basis, and the projective form travel together as the statement “events are pointer-resolved jumps riding on irreversible record formation.” Second, nothing stochastic is added to nature beyond what the bath already supplies:  $\Lambda$  is not a new noise field (contrast GRW/CSL) but the already-derived record-formation rate; the event law assigns ontological weight to monitoring that is physically occurring. Third, the law contains *no* freedom in the probability assignment—that is Theorem 1. Fourth, the law retroactively sharpens the anchoring functional itself: the no-event survival probability is

$$S(t) = \exp\left[-\int_0^t \Lambda ds\right] = e^{-\Phi(t)}, \quad P_{\text{event}}(t) = 1 - e^{-\Phi(t)}, \quad (4)$$

so  $\Phi$  is the cumulative event hazard and the off-diagonal suppression  $e^{-\Phi}$  is the survival probability—one object, two readings.  $\Phi = 1$  is the e-folding scale (63.2% cumulative event probability), not a boundary; the historical threshold language is shorthand, and the threshold-arbitrariness objection is dissolved rather than answered. The Born statistics follow in one line: with populations conserved by the dephasing flow,  $P(k) = \int_0^\infty \lambda_k S dt = p_k \int_0^\infty \Lambda e^{-\Phi} dt = p_k (1 - e^{-\Phi(\infty)})$ , so  $P(k | \text{event}) = p_k$  for any monitoring profile. This formulation supersedes the first-to-threshold race language wherever the two differ.

Stated in standard quantum-trajectory language, the law is an *unraveling*: jump operators  $L_k = \sqrt{\Lambda} P_k$ , no-jump effective Hamiltonian  $H_{\text{eff}} = H - \frac{i\hbar\Lambda}{2}\mathbb{1}$ , jumps  $|\psi\rangle \rightarrow P_k|\psi\rangle/\|P_k|\psi\rangle\|$ . Averaging the trajectories yields exactly one master equation—the dephasing dynamics already present—so anchoring is not a second process layered on decoherence; nothing is double-counted. Three conventions complete the specification. (a) *Repeated jumps*: after a jump into  $P_k$ ,  $p_k = 1$  and further same-channel jumps are identity maps—continuing environmental confirmation, not new events; only state-changing jumps are events. (b) *No event before its record*:  $\Lambda$  counts only irreversible record increments, so an event at any time is accompanied, by construction, by the record formed up to that time; an early event (small  $\Phi$ ) is an event with a small but already-irreversible record. (c) *Scope*: the closed-form Born result holds where the pointer populations are conserved by the flow; in general  $P(k) = \int_0^\infty \Lambda p_k(t) e^{-\Phi} dt$ , and the extension to POVM-smearred continuous outcomes (finite-resolution Kraus operators  $M_x \propto e^{-(\hat{x}-x)^2/4\sigma^2}$  replacing exact projectors, avoiding unbounded momentum injection) is listed open.

### 3 Four Theorems

**Theorem 1** (Born uniqueness from no-signalling). *Assume: (A) standard Hilbert-space kinematics; (B) hazards local in the pointer projectors,  $\lambda_k = \Lambda f(p_k)$  with a single universal nonnegative  $f$ ; (C) mutually exclusive, normalized pointer outcomes; (D) ensemble equivalence—proper and improper mixtures with the same  $\rho$  receive the same hazards; (E) local instruments. Then the unconditioned (event-unobserved) evolution of  $\rho_S$  is affine in  $\rho_S$  if and only if  $f(p) \propto p$ . For any nonlinear  $f$ , an entangled partner’s marginal statistics change when local events fire unobserved—superluminal signalling. Hence within (A)–(E) the linear hazard, and with it the Born weighting, is the unique no-signalling choice. The claim is explicitly not that no-signalling alone generates Hilbert-space probability theory; the kinematics are assumed, the probability rule is derived.*

*Proof.* The unconditioned generator contributed by events is  $\mathcal{L}_{\text{ev}}[\rho] = \Lambda \sum_k f(p_k) \left( \frac{P_k \rho P_k}{p_k} - \rho \right)$ . The term  $\frac{f(p_k)}{p_k} P_k \rho P_k$  is linear in  $\rho$  iff  $f(p)/p$  is constant. For nonlinear  $f$ , prepare  $|\psi\rangle = \sqrt{p}|00\rangle + \sqrt{1-p}|11\rangle$ ; Alice’s unobserved events leave Bob the marginal  $w_0 = f(p)/[f(p) + f(1-p)] \neq p$ . Numerically:  $f(p) = p^2$ ,  $p = 0.7$  shifts Bob’s population from 0.700 to 0.845—an operational signal. For  $f(p) = cp$ ,  $w_0 = p$  identically.  $\square$

**Corollary 1** (Born statistics). *In the measurement regime (monitoring fast relative to system dynamics, so populations are conserved between events—exactly true for pure dephasing), the probability that the first event is  $k$  equals  $p_k(0)$ , for any monitoring profile  $\Lambda(t)$ . Verified semi-analytically:  $p_0 = (0.5, 0.3, 0.2)$  yields  $P(k) = (0.5, 0.3, 0.2)$  under a strongly time-varying  $\Lambda$ .*

This answers the standing objection directly: the Born-shaped hazard is not inserted—it is the unique member of the event class that does not turn entanglement into a telegraph. The same conclusion follows independently from outcome coarse-graining: requiring that merging two outcomes into one coarse outcome commute with the law ( $\lambda_{k_1 \cup k_2} = \lambda_{k_1} + \lambda_{k_2}$  with the same universal  $f$ ) forces Cauchy additivity  $f(x+y) = f(x) + f(y)$ , hence linearity. Two independent routes, one answer. This is consonant with the known linearity constraints on collapse models (Gisin); the contribution here is its statement as a uniqueness theorem internal to ACT’s event class.

**Theorem 2** (No-signalling). *For a bipartite system, Alice’s local event maps satisfy  $\text{Tr}_A \left[ \sum_k P_k^A \rho P_k^A \right] = \text{Tr}_A[\rho]$ . Bob’s reduced state, hence every local expectation value, is invariant under Alice’s unobserved anchoring events.*

**Theorem 3** (Frame-independent joint statistics for commuting spacelike event instruments). *Let Alice’s and Bob’s event maps act on commuting algebras,  $[P_a^A \otimes \mathbb{1}, \mathbb{1} \otimes P_b^B] = 0$ . Then the joint outcome distribution  $P(a, b) = \text{Tr}[(P_a^A \otimes P_b^B) \rho (P_a^A \otimes P_b^B)]$  is identical whether Alice’s event precedes Bob’s or conversely (verified: difference = 0 exactly). Consequently all observable statistics are frame-independent. The model’s internal event sequence for spacelike pairs is not frame-independent; all observable joint statistics are order-independent (Theorem 3), and ACT’s working position is that the internal ordering is unobservable bookkeeping—whether it is ontologically gauge awaits the covariant formulation, listed open. Bell-violating correlations are inherited from the shared pre-anchored state, exactly as the manuscript’s Bell section states, and the price already acknowledged there (nonfactorizable causal structure) is unchanged by the event law. A fully covariant microdynamic remains open.*

**Theorem 4** (Ensemble consistency). *Averaging the PDP over its own stochasticity yields the generator  $\mathcal{L}[\rho] = \Lambda (\sum_k P_k \rho P_k - \rho)$ : the Joos–Zeh pointer-basis dephasing master equation with rate  $\Lambda$ . (Linearity verified numerically; mixture error = 0.) The event law therefore reproduces every ensemble-level prediction of standard open-system quantum mechanics, and is empirically equivalent to standard quantum-trajectory statistics wherever the environment is genuinely monitored. Its new content is ontological—events occur objectively, observer or no observer—plus whatever rate contribution the mass channel adds to  $\Lambda$ .*

## 4 Energy Accounting, Stated Plainly

Per event, the expected post-event energy is  $\sum_k p_k \langle k | H | k \rangle$ , which differs from the pre-event  $\langle H \rangle$  by exactly the coherence energy  $\Delta E_{\text{coh}} = \sum_{j \neq k} \text{Re} \rho_{jk} H_{kj}$  (worked example: offset 0.099 on an  $\mathcal{O}(2)$  Hamiltonian; zero for pointer-diagonal states). Under the flow, this same coherence energy is what the FDT-paired bath exchange is already dissipating; in the measurement regime the offset is bounded by the off-diagonal norm that has, by construction, decayed below  $e^{-\Phi}$  at event time. The honest summary: **energy is conserved in expectation, with per-event fluctuations bounded by the residual coherence energy at threshold and drawn against the bath.** Exact per-event conservation is not claimed; whether the event map can be refined to conserve energy eventwise (e.g., by augmenting  $\mathcal{J}_k$  with a bath-side compensation) is an open problem, listed below.

## 5 What the Law Answers

The six questions every reviewer posed, answered within the model: *What becomes actual?* The pointer component  $k$ , with its local record. *What happens to unrealized components?* Causally inert; they enter no future hazard or coupling; unrealized spectral possibility. *Does the global state change?* Yes—Eq. (3), applied jointly for entangled systems; Theorems 2–3 ensure this is operationally local and frame-consistent. *How are spacelike events coordinated?* Observable joint statistics are order-invariant (Theorem 3); the ontological status of the internal ordering awaits the covariant formulation. *Is no-signalling maintained?* Theorem 2, exactly. *Energy?* Conserved in expectation; per-event accounting bounded and stated.

Wigner’s friend now has a dynamical answer rather than a bare assertion: once the friend’s apparatus crosses threshold, the hazard has fired with probability one at macroscopic  $\Lambda$ , the event map has applied, and Wigner’s unitary description fails *because a jump occurred*—a statement internal to the model, falsifiable with it. The quantum-eraser distinction is built in at Eq. (2):

reversible which-path entanglement contributes zero hazard, so erasure experiments proceed exactly as observed.

## 6 The Ledger

**Postulated (one item):** events are pointer-resolved projective jumps riding on irreversible record formation—the PDP form of the Postulate, including the einselected basis and the identification of  $\Lambda$  with  $d\Phi/dt|_{\text{irr}}$ . This is ACT’s successor to the bridge postulate: sharper, but still a postulate. The claim that *this* unraveling is physically real is addressed by the companion note *Why This Unraveling?* (June 2026): under the Record Condition—events condition only on redundantly recorded, fragment-accessible environmental data—conjugate-conditioned unravelings of every type are excluded, the einselected and recorded observables coincide, and the pointer basis becomes an output of environmental redundancy; within the piecewise-deterministic class the pointer-jump law is then unique, while record-conditioned continuous localization in the pointer manifold remains admissible. What remains irreducibly postulated: the ontic status of record-conditioned jumps, and the discrete-jump form as against such diffusion.

**Derived (four items):** Born weighting, uniquely, from no-signalling (Theorem 1, two independent routes); no-signalling itself (Theorem 2); frame-independence of all observable statistics (Theorem 3); exact consistency with ensemble quantum mechanics and standard trajectory theory (Theorem 4).

**Resolved by the law itself (one item):** the  $\Phi = 1$  conventionality. Equation (4) shows no threshold exists to be conventional:  $\Phi$  is a cumulative hazard and  $\Phi \sim 1$  its e-folding scale.

**Open (three items):** (i) per-event energy conservation, in particular for jump projectors that do not commute with the kinetic energy; (ii) a covariant field-theoretic formulation in which the frame-independence of Theorem 3 is manifest as a property of the ontology rather than verified pairwise for commuting instruments; (iii) the microscopic universal channel and window derivation, unchanged from the Supplement’s constraint section. Also open at the level of generality: the jump map for degenerate, nonprojective, or continuous outcome families, and the non-Markovian regime.

## 7 What This Changes

ACT’s structure before this note: record-formation dynamics + bridge postulate. After: record-formation dynamics + *one sharp stochastic law* from which the Born rule, no-signalling, frame-independent statistics, and ensemble consistency are theorems. That is the completeness standard GRW set for collapse models in 1986—an explicit, falsifiable-in-principle stochastic dynamics—reached here without introducing a fundamental noise field, because ACT’s events ride on environmental monitoring that is already physically present. The manuscripts should adopt the law as “candidate event law” with this note’s ledger verbatim; Lecture 8’s open-steps list shortens by one (the Born hazard is no longer assumed) and gains one (the reality-of-this-unraveling clause now carries the postulate’s weight). The Scientific Status box gains a row: *Derived, conditional on the event class: Born weights, no-signalling, order-invariance, ensemble consistency.*